

# Mechanistic and dual-level direct dynamics studies on the reaction $\text{Cl} + \text{CH}_2\text{FCl}$

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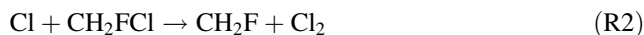
**Abstract** Theoretical investigations are carried out on the reaction  $\text{Cl} + \text{CH}_2\text{FCl}$  by means of direct dynamics method. The minimum energy path (MEP) is obtained at the MP2/6-311G(d, p) level. The energetic information is further improved by single-point energy calculations using QCISD(T)/6-311++G(d, p) method. The kinetics of this reaction are calculated by canonical variational transition state theory incorporating with the small-curvature tunneling correction over a wide temperature range of 220–3,000 K, and rate constant expression are found to be  $k(T) = 1.48 \times 10^{-17} T^{2.04} \exp(-913.91/T)$ . For the title reaction, H-abstraction reaction channel is the major channel at the lower temperatures. At higher temperatures, the contribution of Cl-abstraction reaction channel should be taken into account.

**Keywords** Dual level dynamics ·  $\text{Cl} + \text{CH}_2\text{FCl}$  · Minimum energy path (MEP)

## 1 Introduction

The importance of haloalkanes in atmospheric chemistry is well established [1–5]. Chlorine atom chemistry is of great importance over a wide temperature range. At low temperatures (200–300 K), Chlorine atom plays an active role as ozone destroyer in the atmosphere, and it is important for combustion chemistry at high temperatures (1,000–2,500 K). Hydrochlorofluorocarbons (HCFCs) are emitted into the atmosphere from a number of sources. HCFCs contain Cl atoms and thus contribute to ozone depletion. Chlorine atom reactions play a role in the destruction of these species in the stratosphere [6]. Halogen-substituted hydrocarbons, attacked by Cl atom in the gas phase, also play an important role in the process of industrial chlorination and in the incineration of hazardous halogenated wastes [7]. Experimentally and theoretically, the mechanism and kinetics of such reactions of haloalkanes with various free radicals have been extensively studied and continue to receive considerable attention [8–15].

$\text{CH}_2\text{FCl}$  is a classical species of halogen-substituted hydrocarbons, and its reaction with Cl atoms can proceed via three reaction channels, e.g.



There are some experimental studies concerning rate constants for the title reaction [16–21]. In 1988, Tschuikow-Roux and co-workers [16] reported the H-abstraction

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channel of the title reaction by using the competitive method with methane as a primary standard, and the experimental rate constant was obtained as  $(1.91 \pm 0.57) \times 10^{-13} \text{ cm}^3/(\text{molecule s})$  at 298 K. In 1992, Tuazon et al. [18] studied the rate constant of the title reaction by means of a relative rate technique, and the experimental rate constant was determined as  $(1.1 \pm 0.25) \times 10^{-13} \text{ cm}^3/(\text{molecule s})$  at 298 K. The other experimental rate constant of the reaction  $\text{CH}_2\text{FCl} + \text{Cl} \rightarrow \text{HCl} + \text{CHFCl}$  was determined as  $1.0 \times 10^{-11} e^{[-11.141 \pm 4.124]/RT}$  [in  $\text{cm}^3/(\text{molecule s})$ ] given by Atkinson and co-workers in the temperature region of 273–368 K [21]. The temperature used in the experiment does not cover the whole temperature range of practical interest; therefore, the accurate extrapolation of rate constants to higher temperatures requires the theoretical study. Because no experimental information is available on the branching ratio of the rate constants of the title reaction, theoretical investigation is desirable to give a further understanding of the reaction mechanism and the dynamics property of the multiple channel reaction  $\text{Cl} + \text{CH}_2\text{FCl}$ . To our best knowledge, no previous theoretical work has addressed this reaction.

In this paper, the mechanism of Cl reaction with  $\text{CH}_2\text{FCl}$  is investigated and dual-level direct dynamics studies [22–26] are carried out on this reaction to obtain the rate constants over a wide temperature region of 220–3,000 K. The potential energy surface information, including geometries, energies, gradients, and force constants of the stationary points (reactants, products, complexes and transition states) and extra points along the minimum energy path (MEP), is obtained directly from electronic structure calculations. Subsequently, the rate constants are calculated by the variational transition state theory (VTST) proposed by Truhlar and co-workers [27, 28]. The comparison between theoretical and experimental results has been discussed.

## 2 Computational methods

All the electronic structure calculations were carried out with Gaussian 03 program package [29]. The equilibrium geometries and frequencies of the reactant, products, complexes and transition states are optimized at the restricted or unrestricted second-order perturbation Møller-Plesset level of theory (MP2) [30] using the standard 6-311G(d, p) basis set. At the same level, the MEP is obtained by intrinsic reaction coordinate (IRC) theory in mass-weighted Cartesian coordinates with a gradient step-size  $0.05 (\text{amu})^{1/2} \text{ bohr}$ . Furthermore, the energy profile is refined at the QCISD(T)/6-311++G(d, p) [31] level based on the MP2/6-311G(d, p) geometries. The force constant matrices of the stationary points and selected nonstationary

points near the transition state along the MEP have been also calculated in order to do the following kinetics calculations.

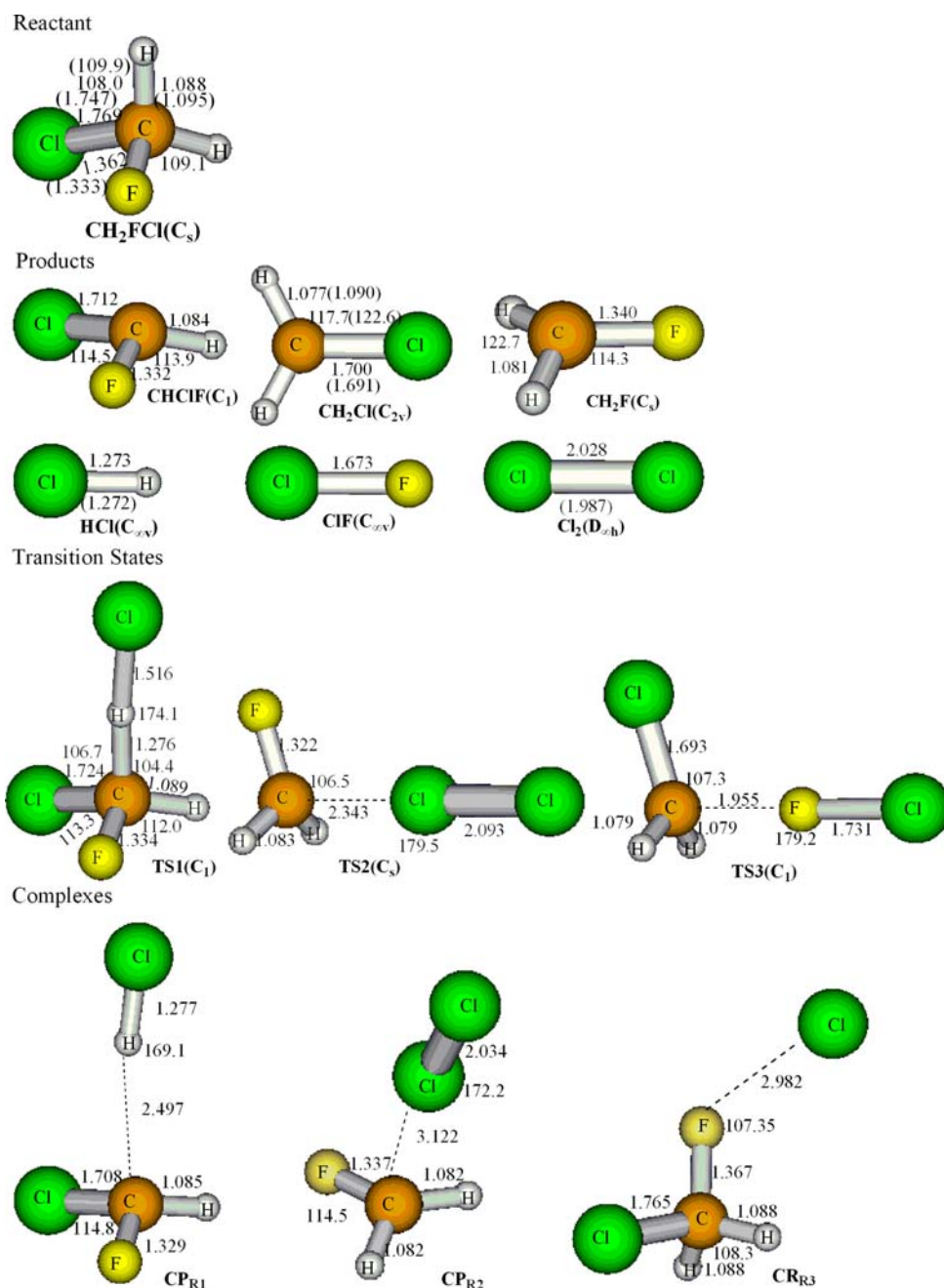
By means of the POLYRATE 9.1 program [32], the dynamics calculations are performed by using the variational transition-state theory (VTST) [27, 28] with the interpolated single-point energies (ISPE) method [33]. The ISPE method is a dual-level direct dynamics scheme that uses a low-level MEP and corrects the energy by interpolating the energy differences at some points along this low-level MEP and single-point energy calculations at a higher-level. The specific form of VTST that we used is canonical variational transition-state theory (CVT) [34–37] with the small-curvature tunneling (SCT) [38, 39] contributions proposed by Truhlar and co-workers. In the calculation of the reactant electronic partition function, two electronic states of Cl atoms,  $^2P_{1/2}$  and  $^2P_{3/2}$ , are included, with an  $881 \text{ cm}^{-1}$  splitting due to spin-orbit coupling. All of the vibrational modes were treated as quantum-mechanical separable harmonic oscillators. During the kinetic calculations, the Euler single-step integrator with a step size of  $0.0001 (\text{amu})^{1/2} \text{ bohr}$  is adopted to follow the MEP, and the generalized normal-mode analysis is performed every  $0.01 (\text{amu})^{1/2} \text{ bohr}$ . The curvature components are calculated using a quadratic fit to obtain the derivative of the gradient with respect to the reaction coordinate.

## 3 Results and discussion

### 3.1 The stationary points

The optimized geometric parameters of the reactant ( $\text{CH}_2\text{FCl}$ ), products ( $\text{CHFCl}$ ,  $\text{HCl}$ ,  $\text{CH}_2\text{F}$ ,  $\text{Cl}_2$ ,  $\text{CH}_2\text{Cl}$ , and  $\text{ClF}$ ), complexes ( $\text{CP}_{R1}$ ,  $\text{CP}_{R2}$ ,  $\text{CR}_{R3}$ ) and transition states (TS1, TS2, and TS3) for three reaction channels, reaction R1, reaction R2, and reaction R3, calculated at the MP2/6-311G(d, p) level along with the available experimental data [40–42] are presented in Fig. 1. From Fig. 1, we can see that the largest deviation between the theoretical bond lengths and the experimental values is  $0.041 \text{ \AA}$  ( $r(\text{Cl}-\text{Cl})$  for  $\text{Cl}_2$ ) and the largest deviation of the bond angles is  $4.8^\circ$  for  $\angle\text{H}_1\text{CCl}$  in  $\text{CH}_2\text{Cl}$ . It is obvious that the theoretical geometric parameters are in good agreement with the experimental data. Furthermore, there are two product complexes ( $\text{CP}_{R1}$ ,  $\text{CP}_{R2}$ ) on the products sides for R1 and R2, and one reactant complex ( $\text{CR}_{R3}$ ) on reactant side in reaction R3. Good agreement between the theoretical and experimental values of  $\text{CH}_2\text{FCl}$ ,  $\text{CH}_2\text{Cl}$ ,  $\text{HCl}$ , and  $\text{Cl}_2$  might imply that the same accuracy could be expected for the calculated geometries of complexes and transition states. At the MP2/6-311G(d, p) level, the  $\text{C}\cdots\text{H}$  bond

**Fig. 1** Optimized geometries of the reactant, products, complexes and transition states at the MP2/6-311G(d, p) level. The values in parentheses are the experimental values (Ref. [40] for  $\text{CH}_2\text{Cl}$ , Ref. [41] for  $\text{HCl}$  and  $\text{Cl}_2$ , Ref. [42] for  $\text{CH}_2\text{FCl}$ ). Bond lengths are in angstroms, and angles are in degrees



distance in  $\text{CP}_{\text{R1}}$  is 2.497 Å, while the other bond lengths are very close to those of the products.  $\text{CP}_{\text{R1}}$  is a weakly hydrogen-bonded complex between the  $\text{HCl}$  molecule and the  $\text{CHFCl}$  radical. The hydrogen bonding occurs between the hydrogen atom of  $\text{HCl}$  and the methyldyne carbon of  $\text{CHFCl}$ . The transition states  $\text{TS1}$  for reaction **R1** and  $\text{TS3}$  for reaction **R3** have the same symmetry, i.e.,  $\text{C}_1$  symmetry. For reaction **R2**, the transition state  $\text{TS2}$  has  $\text{C}_s$  symmetry. In the transition state structure  $\text{TS1}$ , the breaking C–H bond stretches by 18% over the C–H regular bond length in  $\text{CH}_2\text{FCl}$ , and the product-forming H–Cl bond is elongated by about 24% compared with the equilibrium

bond length in isolated  $\text{HCl}$ . Thus, reaction **R1** proceeds via early barrier. In the structure of  $\text{TS2}$  and  $\text{TS3}$ , the C–Cl and C–F bonds that will be broken, increase by 57 and 59% compared with the C–Cl and C–F equilibrium bond lengths in  $\text{CH}_2\text{FCl}$ , and the forming Cl–Cl and F–Cl bonds are about 6 and 5% longer than the regular bond length of the isolated  $\text{Cl}_2$  and  $\text{ClF}$ , respectively. This implies that the two transition states are both near the corresponding products, and these two reactions will both proceed via late transition states.

Table 1 gives the harmonic vibrational frequencies of the reactant, products, complexes, and transition states at

**Table 1** Calculated and experimental frequencies ( $\text{cm}^{-1}$ ) for the reactant, products, complexes and transition states for the title reaction at the MP2/6-311G(d, p) level

Species	MP2/6-311G(d, p)	Expt
CH <sub>2</sub> FCI	398, 799, 1,043, 1,120, 1,296, 1,439, 1,526, 3,147, 3,229	385 <sup>a</sup> , 760 <sup>a</sup> , 1,001 <sup>a</sup> , 1,068 <sup>a</sup> , 1,236 <sup>a</sup> , 1,351 <sup>a</sup> , 1,468 <sup>a</sup> , 2,993 <sup>a</sup> , 3,048 <sup>a</sup>
CHCIF	419, 804, 899, 1,199, 1,361, 3,242	410 <sup>b</sup> , 757 <sup>b</sup> , 910 <sup>b</sup> , 1,250 <sup>b</sup>
HCl	3,090	2,991 <sup>a</sup>
CH <sub>2</sub> F	713, 1,205, 1,212, 1,519, 3,195, 3,353	1,163 <sup>b</sup> , 1,170 <sup>b</sup> , 3,044 <sup>b</sup> , 318 <sup>b</sup>
Cl <sub>2</sub>	539	549 <sup>c</sup>
CH <sub>2</sub> Cl	277, 877, 1,048, 1,473, 3,236, 3,390	389 <sup>b</sup> , 829 <sup>b</sup> , 1,391 <sup>b</sup> , 3,055 <sup>b</sup>
CIF	734	
CP <sub>R1</sub>	19, 35, 64, 208, 216, 422, 819, 904, 1,205, 1,361, 3,031, 3,237	
CP <sub>R2</sub>	17, 39, 55, 95, 153, 539, 744, 1,212, 1,213, 1,519, 3,192, 3,352	
CR <sub>R3</sub>	3, 63, 74, 398, 803, 1,043, 1,104, 1,297, 1,440, 1,524, 3,153, 3,235	
TS1	1204 <i>i</i> , 91, 163, 374, 453, 862, 989, 1,092, 1,180, 1,257, 1,377, 3,184	
TS2	1,309 <i>i</i> , 76, 129, 211, 498, 564, 868, 933, 1,073, 1,483, 3,232, 3,377	
TS3	443 <i>i</i> , 69, 85, 191, 416, 522, 1,005, 1,229, 1,229, 1,519, 3,185, 3,341	

<sup>a</sup> Ref. [43]<sup>b</sup> Ref. [44]<sup>c</sup> Ref. [45]**Table 2** Reaction enthalpies ( $\Delta H_{298}^{\circ}$ ), the relative energies with ZPE correction of complexes ( $\Delta E^c$ ) and potential barrier of the transition states ( $\Delta E^{\text{TS}}$ ) (kcal/mol), classical barrier height ( $V_{\text{MEP}}$ ), vibrationally adiabatic barrier height ( $V_a^G$ ) for the three reactions and together with the experimental values

		MP2/6-311G (d, p)	QCISD(T)/6-311++ G(d, p)//MP2/6-311G(d, p)	Expt <sup>a</sup>	
CH <sub>2</sub> FCI + Cl → CHFCI + HCl	$\Delta H_{298}^{\circ}$	-4.83	-1.13	-3.0 ± 2.4	
	$\Delta E^{\text{TS}} + \text{ZPE}$	5.02	3.64		
	CP <sub>R1</sub>	$\Delta E^c + \text{ZPE}$	-2.26		-2.73
		$V_{\text{MEP}}$	9.27		7.89
		$V_a^G$	25.04		23.65
	CH <sub>2</sub> FCI + Cl → CH <sub>2</sub> F + Cl <sub>2</sub>	$\Delta H_{298}^{\circ}$	37.85		37.29
$\Delta E^{\text{TS}} + \text{ZPE}$		43.39	37.73		
CP <sub>R2</sub>		$\Delta E^c + \text{ZPE}$	39.20	35.90	
		$V_{\text{MEP}}$	45.12	39.45	
		$V_a^G$	63.41	57.74	
CH <sub>2</sub> FCI + Cl → CH <sub>2</sub> Cl + CIF		$\Delta H_{298}^{\circ}$	55.89	55.45	49.9 ± 2.0
	$\Delta E^{\text{TS}} + \text{ZPE}$	76.29	62.89		
	CR <sub>R3</sub>	$\Delta E^c + \text{ZPE}$	-1.30	-1.96	
		$V_{\text{MEP}}$	78.52	65.11	
		$V_a^G$	96.31	82.90	

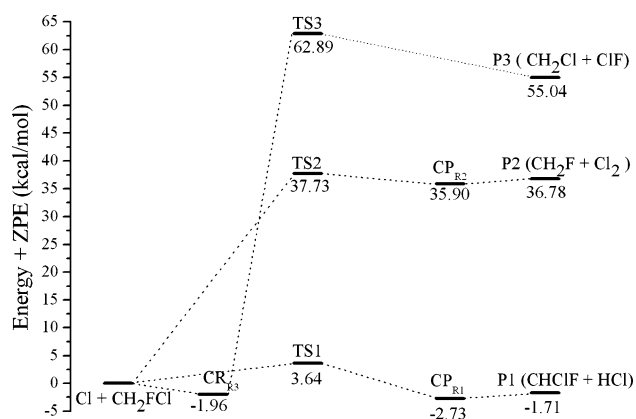
<sup>a</sup> Experimental value derived from the standard heats of formation (in kcal/mol): [46] for CH<sub>2</sub>FCI (-62.6 kcal/mol); HCl (-22.1 kcal/mol); Cl (29.0 kcal/mol); CIF (-12.0 kcal/mol); [47] for CHFCI (14.51 ± 2.4 kcal/mol) [48] for CH<sub>2</sub>Cl (28.3 ± 2.0 kcal/mol)

the MP2/6-311G(d, p) level as well as the corresponding experimental results [43–45]. For the species of reactant and products, the calculated frequencies are in good agreement with the experimental values, for each transition state yields one and only one imaginary frequency corresponding to the stretching modes of the coupling breaking and forming bonds. The values of those imaginary frequencies are 1204.0*i*, 1309.3*i*, and 443.5*i*  $\text{cm}^{-1}$  for reaction R1, R2, and R3, respectively.

The reaction enthalpies ( $\Delta H_{298}^{\circ}$ ), the relative energies of hydrogen-bonded complexes ( $\Delta E^c$ ), and the potential barrier of the transition states ( $\Delta E^{\text{TS}}$ ) (kcal/mol)

with zero-point (ZPE) corrections for reactions R1, R2, and R3 calculated at the MP2/6-311G(d, p) level and QCISD(T)/6-311++G(d, p)//MP2/6-311G(d, p) level are listed in Table 2, as well as the available experimental reaction enthalpies [46–48]. The theoretical values at 298 K of  $\Delta H_{298}^{\circ}$ , -1.13 and 55.45 kcal/mol for reactions R1 and R3 are in reasonable agreement with the corresponding experimental values, -3.0 ± 2.4 and 49.9 ± 2.0 kcal/mol.

The profile of potential energy surface of the title reaction obtained at the QCISD(T)/6-311++G(d, p)//MP2/6-311G(d, p)+ZPE level is described in Fig. 2. Note

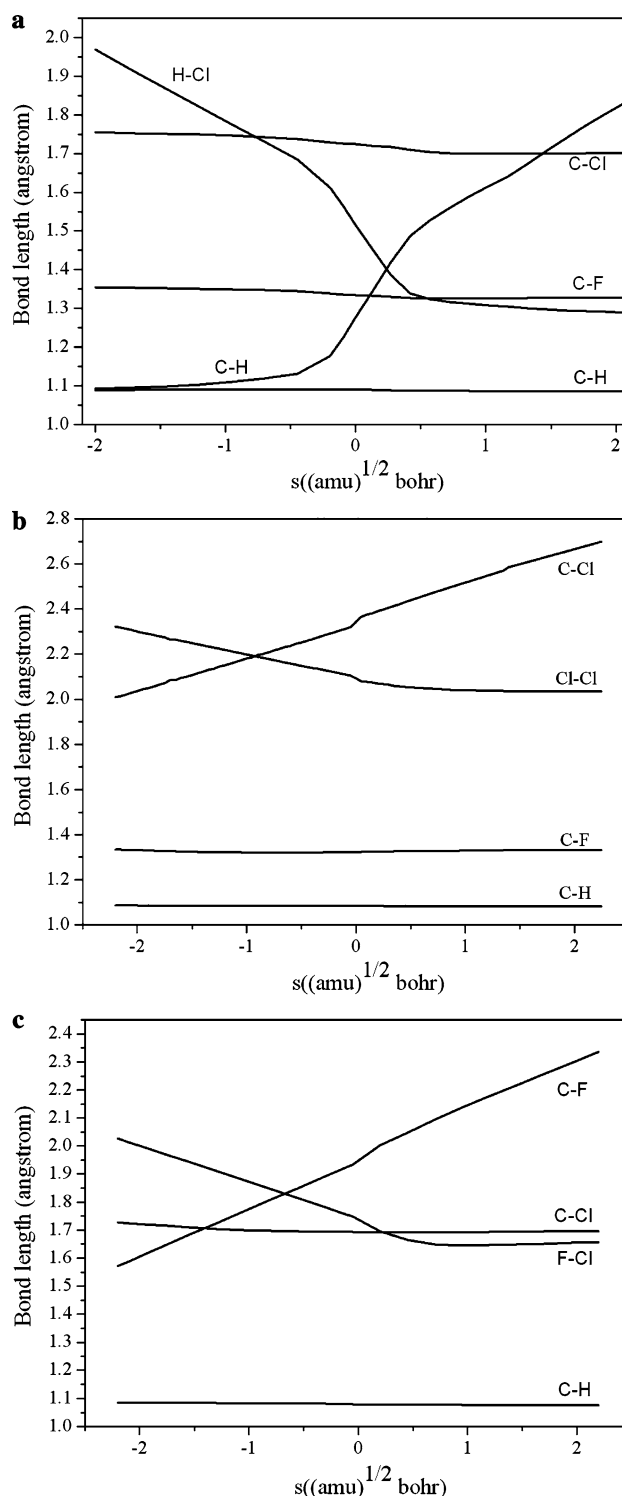


**Fig. 2** Schematic potential energy surface for the title reaction. Relative energies (in kcal/mol) are calculated at the QCISD(T)/6-311++G(d, p)//MP2/6-311G(d, p) + ZPE level

that the energies of reactants are set to be zero for reference. For reaction **R1**, the attack of Cl atom on the C–H bond of  $\text{CH}_2\text{FCl}$  would proceed via a hydrogen-bonded complex ( $\text{CP}_{\text{R}1}$ ) on the product side, which is 2.73 kcal/mol lower than the reactants. For reaction **R2** and **R3**, the product complex ( $\text{CP}_{\text{R}2}$ ) and reactant complex ( $\text{CR}_{\text{R}3}$ ) are located at the exit and the entrance of the channels, respectively, and so they proceed via indirect mechanisms. The energy of complex  $\text{CP}_{\text{R}2}$  is 35.90 kcal/mol higher than that of reactants, and the energy of complex  $\text{CR}_{\text{R}3}$  is 1.96 kcal/mol lower than that of reactants. The barrier heights of reactions **R2** and **R3** are 34.09 and 59.25 kcal/mol higher than that of reaction **R1** (3.64 kcal/mol). Thus, the H-abstraction channel is more favorable, and the halogen abstraction channels are minor pathway. At the same time, the H-abstraction reaction **R1** is exothermic by 1.71 kcal/mol, while reaction **R2** and **R3** are endothermic by 36.78 and 55.04 kcal/mol, respectively, which indicates that reaction **R1** is thermodynamically more favorable than reactions **R2** and **R3**, and the rate constants of reaction **R1** will be larger than those of the other two reactions (**R2** and **R3**). This view will be further testified in the following study of the rate constants.

### 3.2 Minimum energy path

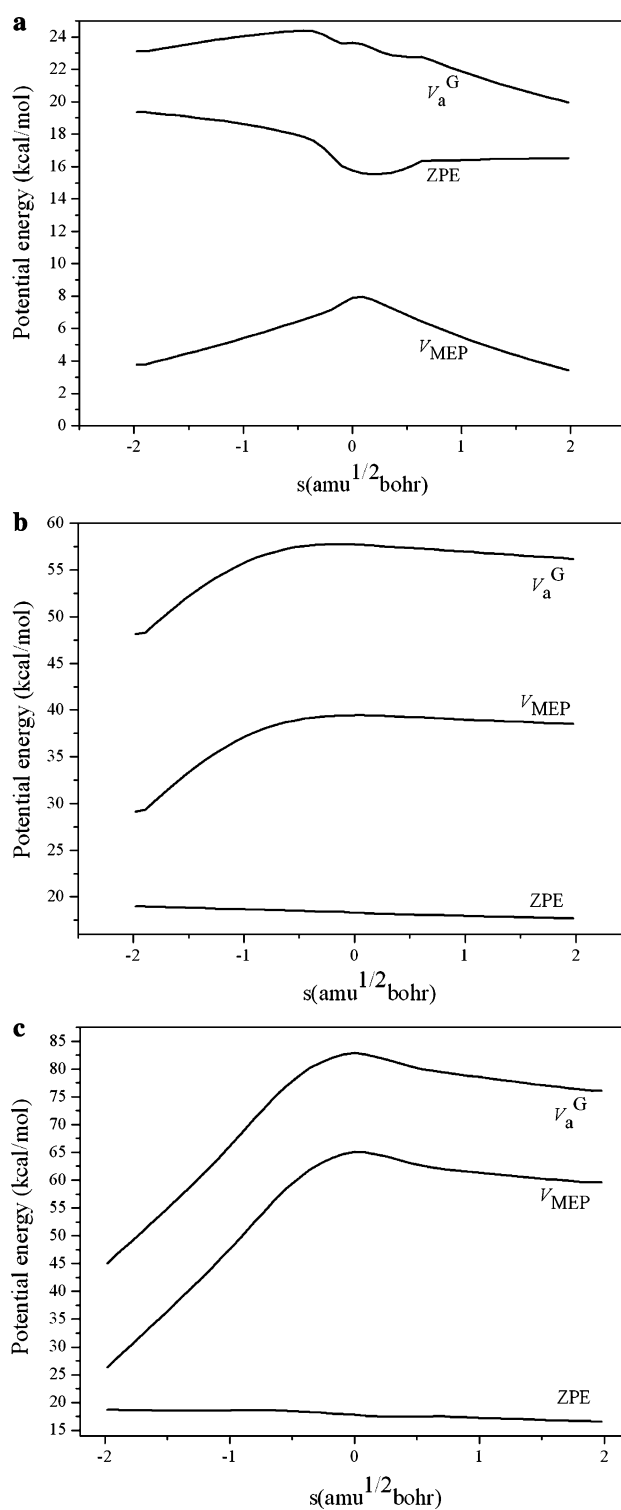
The changes of bond lengths along the IRC for the Cl reaction with  $\text{CH}_2\text{FCl}$  are plotted in Fig. 3. From Fig. 3a, it is easily seen that with the proceeding of the H– abstraction reaction channel, the breaking C–H bond and forming H–Cl bond lengths change smoothly up to the about  $s = -0.45 \text{ (amu)}^{1/2} \text{ bohr}$ . After that, both bonds change



**Fig. 3** **a** Changes of the bond lengths along the MP2/6-311G(d, p) minimum energy path for the  $\text{Cl} + \text{CH}_2\text{FCl} \rightarrow \text{CHCIF} + \text{HCl}$ . **b** Same types of curves as those in **a** for the  $\text{Cl} + \text{CH}_2\text{FCl} \rightarrow \text{CH}_2\text{F} + \text{Cl}_2$ . **c** Same types of curves as those in **a** for the  $\text{Cl} + \text{CH}_2\text{FCl} \rightarrow \text{CH}_2\text{Cl} + \text{ClF}$

rapidly up to about  $s = 0.42 \text{ (amu)}^{1/2} \text{ bohr}$ , and later, the changes become slow again. However, the lengths of the other bond show nearly no change during the entire reaction processes. The points where these remarkable changes take place indicate the beginning of the dissociation of the C–H bond [ $s = -0.45 \text{ (amu)}^{1/2} \text{ bohr}$ ] and the end of the formation of the H–Cl bond [ $s = 0.42 \text{ (amu)}^{1/2} \text{ bohr}$ ]. It is thus evident that hydrogen abstraction from C to Cl essentially takes place in the region of  $s = -0.45$ – $0.42 \text{ (amu)}^{1/2} \text{ bohr}$ . For  $\text{Cl} + \text{CH}_2\text{FCl} \rightarrow \text{CH}_2\text{F} + \text{Cl}_2$  reaction channel (see Fig. 3b), the breaking C–Cl bond the forming Cl–Cl bond lengths change linearly along the reaction coordinate up to about  $s = -0.05 \text{ (amu)}^{1/2} \text{ bohr}$ . From  $s = -0.05$ – $0.05 \text{ (amu)}^{1/2} \text{ bohr}$  both bonds change rapidly. After that, the breaking C–Cl bond length increases linearly again, and the forming Cl–Cl bond length decreases slowly until the reaction coordinate  $s$  reaches  $0.70 \text{ (amu)}^{1/2} \text{ bohr}$  and remains unchanged thereafter. This indicates that the beginning of the formation of the Cl–Cl bond is at  $s = -0.05 \text{ (amu)}^{1/2} \text{ bohr}$  and the end of the formation of the Cl–Cl bond is at  $s = 0.70 \text{ (amu)}^{1/2} \text{ bohr}$ . While the other bond lengths show nearly no change during the entire reaction processes. For the  $\text{Cl} + \text{CH}_2\text{FCl} \rightarrow \text{CH}_2\text{Cl} + \text{ClF}$  reaction channel (see Fig. 3c), the bond length changes along the IRC are similar to the reaction channel  $\text{Cl} + \text{CH}_2\text{FCl} \rightarrow \text{CH}_2\text{F} + \text{Cl}_2$ . The different results between Fig. 3a and Fig. 3b or c are due to the structure of the transition state. For reaction R1, it proceeds via early barrier, while for reaction R2 and R3, they both proceed via late transition states.

The minimum-energy paths (MEPs) of the three reaction channels are calculated by the intrinsic reaction coordinate (IRC) theory at the MP2/6-311G(d, p) level. Using interpolated single-point energies (ISPE) approach, the potential profiles are further refined by four single-point calculations (ISPE-4) at QCISD(T)/6-311++G(d, p)/MP2/6-311G(d, p) level. Figure 4a–c depicts the classical potential energy curves  $V_{\text{MEP}}(s)$  and the ground-state vibrational adiabatic potential energy curves  $V_{\text{a}}^{\text{G}}(s)$  as functions of the intrinsic reaction coordinate  $s$  at QCISD(T)/MP2 level together with the zero-point energy curves  $\text{ZPE}(s)$ , where  $V_{\text{a}}^{\text{G}}(s) = V_{\text{MEP}}(s) + \text{ZPE}(s)$ , for the three reaction channels. Table 2 lists the classical barrier height ( $V_{\text{MEP}}$ ) and the vibrationally adiabatic ground state barrier height ( $V_{\text{a}}^{\text{G}}$ ). In Fig. 4a, it can be seen that the curve of  $V_{\text{a}}^{\text{G}}(s)$  is different in shape from the corresponding curve of  $V_{\text{MEP}}(s)$ . We can interestingly find that for H-abstraction reaction channel, the ground-state vibrational adiabatic potential surface  $V_{\text{a}}^{\text{G}}(s)$  has a local maximum at  $s = -0.45 \text{ (amu)}^{1/2} \text{ bohr}$  in the reactant side. The nonregular  $V_{\text{a}}^{\text{G}}(s)$  shape may be attributed to the combination of two different factors: the low-energy



**Fig. 4** a Classical potential energy ( $V_{\text{MEP}}(s)$ ), zero-point energies (ZPE), and vibrational adiabatic potential energy ( $V_{\text{a}}^{\text{G}}(s)$ ) as a function of the reaction coordinate,  $s$ , at the QCISD(T)/MP2 level for the  $\text{Cl} + \text{CH}_2\text{FCl} \rightarrow \text{CH}_2\text{F} + \text{Cl}_2$ . b Same types of curves as those in a for the  $\text{Cl} + \text{CH}_2\text{FCl} \rightarrow \text{CH}_2\text{F} + \text{Cl}_2$ . c Same types of curves as those in a for the  $\text{Cl} + \text{CH}_2\text{FCl} \rightarrow \text{CH}_2\text{Cl} + \text{ClF}$

barrier and the relatively significant early drop in the zero-point energy (ZPE) curve prior to the saddle point zone [48]. This is not only a feature of the present reaction systems, but also has been observed in other reactions, as reported by Rosenman et al. [49] and Li et al. [50]. Since the curve of  $V_a^G(s)$  differs obviously from the curve of  $V_{\text{MEP}}(s)$  in shape, the variational effect might be very important in evaluating the rate constants. In order to understand variational effect further, the dynamics bottleneck property of the reaction R1 based on the canonical variational transition state approach are calculated. The dynamics bottleneck properties denote that the position  $s$  of the variational transition state at various temperatures deviates from the saddle point at  $s = 0$ . For the reaction channel R1  $\text{Cl} + \text{CH}_2\text{FCl} \rightarrow \text{CHFCl} + \text{HCl}$ , the maximum value of shifted  $s$  is  $-0.04$  ( $\text{amu}^{1/2}$  bohr) at 220 K, where the values of  $V_{\text{MEP}}(s = -0.04)$  and  $V_a^G(s = -0.04)$  are 9.24 and 25.08 kcal/mol. The differences,  $V_{\text{MEP}}(s = -0.04) - V_{\text{MEP}}(s = 0) = -0.028$  kcal/mol, and  $V_a^G(s = -0.04) - V_a^G(s = 0) = 0.05$  kcal/mol, are small. It can, therefore, be concluded that the variational effect on the rate constant is small. For the remaining two reaction channels, i.e., (R2)  $\text{Cl} + \text{CH}_2\text{FCl} \rightarrow \text{CH}_2\text{F} + \text{Cl}_2$  and (R3)  $\text{CH}_2\text{FCl} \rightarrow \text{CH}_2\text{Cl} + \text{ClF}$  (see Fig. 4b, c), the  $V_{\text{MEP}}(s)$  and  $V_a^G(s)$  curves are similar in shape, and their maximum positions are almost the same. The zero-point energy ZPE curves are practically constants as  $s$  varies with only a gentle fall near the saddle point ( $s = 0$ ). As a result, for the two reaction channels (R2 and R3), the variational effect on the calculation of rate constant will also be small, as will be confirmed by the rate constant calculations.

### 3.3 Rate constants

The rate constants of the individual channel,  $k_1$ ,  $k_2$ , and  $k_3$  are calculated by conventional transition state theory (TST), canonical variational transition state theory (CVT), and the CVT with the small-curvature tunneling (SCT) contributions in a wide temperature range from 220 to 3,000 K at the QCISD(T)/6-311++G(d, p)//MP2/6-311G(d, p) level of theory. The calculated rate constants of the three reaction channels are displayed in Table 3 and plotted against the reciprocal of temperature in Fig. 5a–c together with the corresponding experimental data. From Table 3 and Fig. 5a–c, we can see that for reaction R1, R2, and R3, the CVT rate constants and the TST rate constants are nearly the same over the whole temperature range 220–3,000 K. This means that the variational effect is small for the three reaction channels. For the reaction R1 and R3 channels, the tunneling effect, that is, the ratio between CVT/SCT and CVT rate constant, plays an important role in the calculation of rate constants at low temperatures and becomes negligible in the high temperature range. For example, the ratios of  $k_1(\text{CVT/SCT})/k_1(\text{CVT})$  are 1.71 at

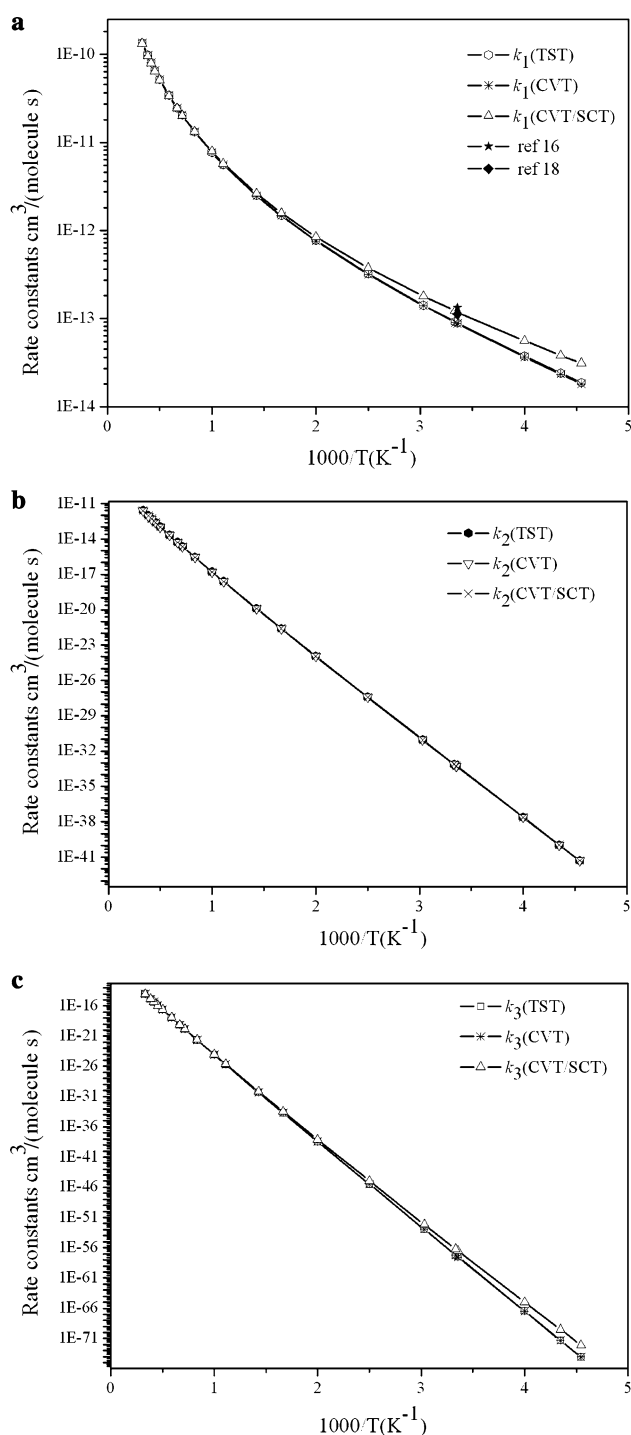
**Table 3** The calculated CVT/SCT rate constants [ $\text{cm}^3/(\text{molecule s})$ ] of three reaction channels and the overall rate constants  $k$  [ $\text{cm}^3/(\text{molecule s})$ ] in the temperature range 220–3,000 K

$T$ (K)	$k_1$	$k_2$	$k_3$	$k$
220	3.12E–14	4.94E–42	8.10E–73	3.12E–14
230	3.84E–14	1.05E–40	2.94E–70	3.84E–14
250	5.60E–14	2.28E–38	1.00E–65	5.60E–14
298	1.19E–13	5.24E–34	3.13E–57	1.19E–13
300	1.22E–13	7.44E–34	6.22E–57	1.22E–13
330	1.80E–13	8.64E–32	6.92E–53	1.80E–13
400	3.78E–13	3.76E–28	1.01E–45	3.78E–13
500	8.47E–13	1.11E–24	7.34E–39	8.47E–13
600	1.58E–12	2.46E–22	3.19E–34	1.58E–12
700	2.62E–12	1.22E–20	7.15E–31	2.62E–12
900	5.74E–12	2.40E–18	2.39E–26	5.74E–12
1,000	7.90E–12	1.57E–17	9.58E–25	7.90E–12
1,200	1.32E–11	2.74E–16	2.55E–22	1.32E–11
1,400	2.02E–11	2.19E–15	1.45E–20	2.02E–11
1,500	2.43E–11	5.09E–15	7.37E–20	2.43E–11
1,700	3.38E–11	2.07E–14	1.10E–18	3.38E–11
2,000	5.10E–11	1.04E–13	2.39E–17	5.11E–11
2,200	6.45E–11	2.43E–13	1.19E–16	6.47E–11
2,400	7.95E–11	5.00E–13	4.57E–16	8.00E–11
2,600	9.60E–11	9.27E–13	1.45E–15	9.69E–11
3,000	1.33E–10	2.55E–12	9.35E–15	1.36E–10

220, 1.35 at 300, 1.19 at 400, 1.08 at 600, 1.03 at 1,000, and 0.99 at 3,000 K, respectively. The ratios of  $k_3(\text{CVT/SCT})/k_3(\text{CVT})$  are 96.10, 10.10, 3.42, 2.15, 1.25, 1.05, and 1.02 at 220, 300, 400, 500, 800, 2,000, and 3,000 K, respectively. For the reaction R2 channel, Fig. 5b shows that the CVT and CVT/SCT rate constants are nearly the same over the whole temperature range 220–3,000 K, and this means that the tunneling effect is small.

The CVT/SCT rate constants of  $k_1$  at 298 K,  $1.19 \times 10^{-13} \text{ cm}^3/(\text{molecule s})$ , are in good agreement with the available experimental value of  $(1.1 \pm 0.25) \times 10^{-13} \text{ cm}^3/(\text{molecule s})$  [18]. The overall theoretical CVT/SCT rate constants  $k$  for the title reaction are calculated from the sum of the corresponding individual rate constants, that is,  $k = k_1 + k_2 + k_3$ . The deviation between the theoretical and experimental values remains within a factor of approximately 1.06. Thus, the present calculations may provide reliable prediction of the rate constants for the title reaction over a wide temperature range.

Table 3 shows that the rate constants of reaction R1,  $k_1$ , are about 27–2 and 58–6 orders of magnitude higher than those of the rate constants of  $k_2$  and  $k_3$ , from 220 to 2,000 K, respectively. Thus, we can conclude that the H-abstraction channel, reaction R1, is always absolute



**Fig. 5** The TST, CVT, and CVT/SCT rate constants as a function of  $10^3/T$  between 220 and 3,000 K for the **a**  $\text{Cl} + \text{CH}_2\text{FCI} \rightarrow \text{CHFCI} + \text{HCl}$ . **b**  $\text{Cl} + \text{CH}_2\text{FCI} \rightarrow \text{CH}_2\text{F} + \text{Cl}_2$ . **c**  $\text{Cl} + \text{CH}_2\text{FCI} \rightarrow \text{CH}_2\text{Cl} + \text{ClF}$ , together with the available experimental values

dominant in the temperature of 220–2,000 K and reactions **R2** and **R3** are minor pathways. It is consistent with the analysis result from the potential barrier heights of these three reaction channels. The contribution of  $k_2$  and  $k_3$  to the overall rate constants increases as the temperature

increases, while reaction **R2** becomes more and more competitive with the temperature increase from 2,000 to 3,000 K.

Finally, for convenience of future experimental measurements, we give three-parameter rate constant expressions by fitting the theoretical results for the three reaction channels and the whole reaction in the temperature range of 220–3,000 K, and the expressions are given as follows: [in unit of  $\text{cm}^3/(\text{molecule s})$ ]

$$k_1(T) = 4.79 \times 10^{-17} T^{1.90} \exp(-1128.81/T),$$

$$k_2(T) = 3.11 \times 10^{-15} T^{1.48} \exp(-15579.36/T),$$

$$k_3(T) = 2.08 \times 10^{-14} T^{1.26} \exp(-32813.08/T),$$

$$k(T) = 1.48 \times 10^{-17} T^{2.04} \exp(-913.91/T).$$

## 4 Conclusion

In this paper, the reaction  $\text{Cl} + \text{CH}_2\text{FCI} \rightarrow \text{products}$  have been studied by dual-level direct dynamics methods. The potential energy surface information is obtained at the MP2/6-311G(d, p) level, and higher-level energies are improved at the QCISD(T)/6-311++G(d, p) level. Three reaction channels are identified, i.e. H-abstraction, Cl-abstraction, and F-abstraction. The calculated potential barriers show that the major pathway is H-abstraction channel leading to the products,  $\text{CHFCI} + \text{HCl}$ . For each individual reaction channel, the theoretical rate constants in the temperature region of 220–3,000 K are calculated by the canonical variational transition state theory (CVT) with the small-curvature tunneling correction (SCT). The calculated total rate constants of these reactions are in good agreement with the corresponding experimental values. The results show that H-abstraction channel is important at lower temperatures; the Cl-abstraction should be considered when the temperature is higher the 2,000 K. The F-abstraction reaction channel is a minor channel over the whole temperature region. The rate constant calculations show that the overall rate constants have positive temperature dependence. It is found that the variational effect is small for the three reaction channels. The small-curvature tunneling effect plays an important role at the lower temperatures for reaction channels **R1** and **R3**, while it is small for the reaction channel **R2**. Our results may be useful for further experimental studies on the dynamical properties of the title reaction.

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